The theoretical view of high-z Clusters

Nelson Padilla, PUC, Chile Pucón, November 2009

The Plan:

I) To see what the observations are telling us using models that agree with the cosmology, and with other observations at different redshifts (as many observations as possible).

II) Cosmological origin of clusters

How does the gas end up heated, cooled, reheated, in stars, in galaxies (how many, how extended), why are there stars outside galaxies, when did all this happen, etc.

This talk:

- Tools to study clusters from the theory side.
- Subject outline
 - Density peaks: when do Clusters form?
 - Physical processes in Clusters: baryon heating, cooling, ...
 - Galaxy formation

Cosmology with Clusters: see Gus Evrard's Talk

Tool I: Linear Theory + analytic approximations Tool II: DM simulations

Density fields

- small initial fluctuations
- Growth via gravitational instability
- Downsizing

$$\delta_{\text{obs}} = b \left(\delta_m - 3 \, \delta z \right) + A + 2D + \left(v^\alpha - B^\alpha \right) e_\alpha + E_{\alpha\beta} e^\alpha e^\beta - \left(1 + z \right) \frac{\partial}{\partial z} \, \delta z - 2 \, \frac{1+z}{Hr} \, \delta z - \delta z - 5p \, \delta \mathcal{D}_L - 2 \, \kappa + \, \frac{1+z}{H} \frac{dH}{dz} \, \delta z + 2 \, \frac{\delta r}{r} \,, \qquad (36)$$

Yoo, Fitzpatrick & Zaldarriaga (2009)

Springel et al. (2005) 1 1 1 1 1 11 10-1 z = 0.00z = 1.50 10-2 z = 3.06 $M^2/p \, dn/dM$ z = 5.7210⁻³ z = 10.07 104 10* 10¹⁰ 10¹¹ 10¹² 10¹³ 10¹⁴ 10¹⁵ 10¹⁶ $M [h^{-1} M_{\odot}]$ Solid: Jenkins et al. Dotted: PS

PressSchechter

Bond 1991 Excursion set approach

Non spherical collapse by Sheth, Mo & Tormen

Main problem: three axes collapse at different epochs, so which is the appropriate δ_c ?

Linear theory for ellipsoids predicts dependence on ellipticity, prolacity and even δ_{c}

Best option for this behaviour:

$$\frac{\delta_{\rm ec}(e,p)}{\delta_{\rm sc}} = 1 + \beta \left[5 \left(e^2 \pm p^2 \right) \frac{\delta_{\rm ec}^2(e,p)}{\delta_{\rm sc}^2} \right]^{\gamma}$$

with β =0.47, γ =0,615 and δ_{sc} =1.7 (sph. collapse)

Typical values for gaussian field, p=0 $e=(\sigma/\delta)/5^{1/2}$

so that,

$$\delta_{
m ec}(\sigma,z) = \delta_{
m sc}(z) \, \left(1 + eta \left[rac{\sigma^2}{\sigma_*^2(z)}
ight]^\gamma
ight)$$



Non spherical collapse by Sheth, Mo & Tormen



Applied to Clusters

Mass function constraints on the Cosmology

• Eke, Cole & Frenk (1996) $\sigma_8 = (0.50 \pm 0.04)\Omega_0^{-0.53 + 0.13\Omega_0}$ for $\Omega_0 + \Lambda_0 = 1.000$

• Sánchez, Padilla & Lambas (1998)



Extended
 Press
 Schechter

$$dN = \frac{1}{\sqrt{2\pi}} \frac{\Delta\omega}{(\Delta\sigma^2)^{3/2}} \exp\left[-\frac{(\Delta\omega)^2}{2\Delta\sigma^2}\right] \left|\frac{d\sigma^2}{dM}\right| \frac{M_0}{M} dM$$
$$\omega = \delta_c / D(t) \Delta\omega = \omega - \omega_0 \Delta\sigma^2 = \sigma^2(M) - \sigma^2(M_0)$$

Bower et al. (1991), Lacey & Cole (1993)

z = 48.4

T = 0.05 Gyr

Acquarius project

Extended Press Schechter

$$rac{dN}{dM_1}
ightarrow rac{dN}{dM_1} \ G(\sigma_1/\sigma_2, \delta_2/\sigma_2)$$

$$G(\sigma_1/\sigma_2, \delta_2/\sigma_2) = G_0 \left(\frac{\sigma_1}{\sigma_2}\right)^{\gamma_1} \left(\frac{\delta_2}{\sigma_2}\right)^{\gamma_2}$$

$$G_0 = 0.61, \gamma_1 = 0.27, \gamma_2 = 0.0$$

Mass distributions of the largest and second largest progenitors



& Helly (2007) Cole Parkinson,

Applied to Clusters

Extended PS and the formation of the first haloes

• Angulo & White (2009)

DM particle: neutralino Fluctuations that survive: 0.7pc or 10⁻⁸M_{sun} and up Affected haloes: much higher masses (do they accrete clumps or difuse matter?) DM simulations? only in 2050. So, extended PS.



Applied to Clusters

Extended PS and the formation of the first haloes • Angulo & White (2009)

Mass of the first collapsed object:

5-7 orders of mag. larger than M_{fs}

Almost flat distribution

Likely tidal disruption of M_{fs} clusters at infall.





Applied to Clusters •Li, Mo & Gao, 2008 (and many other works)

Simple models where the clustering depends only on mass need to be adapted to this phenomenology.



Figure 4. Age dependence of halo bias. Formation time used is indicated in each panel. Dashed lines are for oldest 20% halos while solid lines are for youngest 20% halos; the thick gray lines represent the bias of all the halos regardless of their ages. Error barshows the Poisson error.

There is a definite need for numerical simulations to ensure merger trees respond to the assembly bias.

Tool III: SAMs and Hydro (with stellar population synthesis models)

• Ceccarelli et al 2010 (Infalls via filaments)

• Pivato, Padilla & Lambas (2006)

 Filaments in cosmological numerical simulations (Colberg et al., 2005, Gonzalez & Padilla 2010)

- Ceccarelli et al 2010 (Infalls via filaments)
- Pivato, Padilla & Lambas (2006)



The material arrives preprocessed?

Louise Edwards talk, McGee et al. (2009), Porter et al. (2008), Wilman et al. (2008)

Galaxies evolve within groups before they fall onto Clusters of galaxies.



González & Padilla (2009)

Applied to

Clusters

Similar results in Lagos, Cora & Padilla (2008) SAM

• Filaments in cosmological numerical simulations

Colberg et al. (2005)



González & Padilla (2010)



Sub-halo (galaxy) mergers



Sub-halo (galaxy) mergers







if we have the wrong dynamical friction timescales



Applied to Clusters

Mergers of subhaloes: how many individual peaks in clusters?

- Halo model approaches
- Extent of sub-peaks in clusters: QbC

Talk by Sebastián López Poster by Heather Andrews

Baryonic Mass Infall

- Gas in DM haloes is heated to T_{vir}
- Gas Cooling from Hot Gas reservoir (Rees & Ostriker, 1977, White & Rees, 1978)
- Cool gas via filaments (Binney 1977, Dekel & Birnboim, 2003, Katz et al., 2003)

Baryonic Mass Infall

Gas in DM haloes is heated to Tvir

$$kT_{\rm x} = (1.38 \text{ keV})\beta^{-1}$$
$$\times \left(\frac{M_{\rm vir}}{10^{15} h^{-1} M_{\odot}}\right)^{2/3} \left[\frac{\Omega_0}{\Omega(z)}\right]^{1/3} \Delta_{\rm vir}^{1/3} (1+z),$$

Eke et al. (1996) Voit & Donahue (1998)



Baryonic mass infall halting mechanisms

Gas cooling from hot reservoir

Lagrangian simulations by Birnboim & Dekel (2004)

Birnboim et al., (2007)



Baryonic mass infall halting mechanisms

 Infall of gas via filaments





Dekel et al., 2008

Baryonic mass infall halting mechanisms

• Gas heating and cooling



Star formation and mass infall halting

Forming stars in discs of cooled gas, and associating to this formation SN explosions (enrichment) and AGN.

Interactions in clusters (for small galaxies)

An for massive galaxies, AGN Feedback (Croton et al. 2006, Bower et al., 2006, Lagos, Cora & Padilla, 2008) SN superwinds (Baugh et al., 2005), Gravitational Heating (Birnboim) AIMs: to produce sensible metallicities, old BCGs and Red Sequences

Applied to Clusters

The red sequence from AGN quenching:



Lagos, Cora & Padilla (2008)

Applied to Clusters

The metallicity gradients in clusters using SAMs and Hydro together.

Galaxies need to expel their metals out to distances of 100kpc in order to fit the observations



Cora, Tornatore, Tozzi & Dolag (2008)

Applied to And how do these Clusters galaxies look like today?



Tool IV: stellar population synthesis models

Growth of the red sequence

Allows to interpret results but does not ensure the descendants will look like measured z=0 galaxies.



Do these galaxies evolve into a SDSS LF?

De Lucia et al., 2006 Also talks by Diego Capozzi and by Nicola Menzi (in a couple of minutes)

Sumary

- Tool I: provides direct relation between cosmological model and distributions of haloes. Differences with fully non-linear models help improve the analytic approach.
- Tool II: simulations; these are essential in understanding the different processes ignored in Tool I.
- Tool III: the baryon physics can be followed using different levels of resolution and detail.

These can be used to study high redshift Clusters and at the same time ensuring a reasonable population of final z=0 gals.

• Tool IV: helps understand observations, but on itself does not allow to look after the z=0 population.