

# Supporting the Obscuring Torus by Radiation Pressure

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# The Basic Problem

$N(\text{type 2}) \sim N(\text{type 1})$  implies  $\Delta\Omega_{\text{obscured}} \sim \Delta\Omega_{\text{unobscured}}$

which in turn implies  $h_{\text{obscured}} \sim r$

but  $h/r \sim \Delta v/v_{\text{orb}}$ , yet  $v_{\text{orb}} \sim 100 \text{ km/s}$

If  $\Delta v = c_s$ , then  $T \sim 10^5 \text{ K} \gg T_{\text{subl}} (\text{dust})$

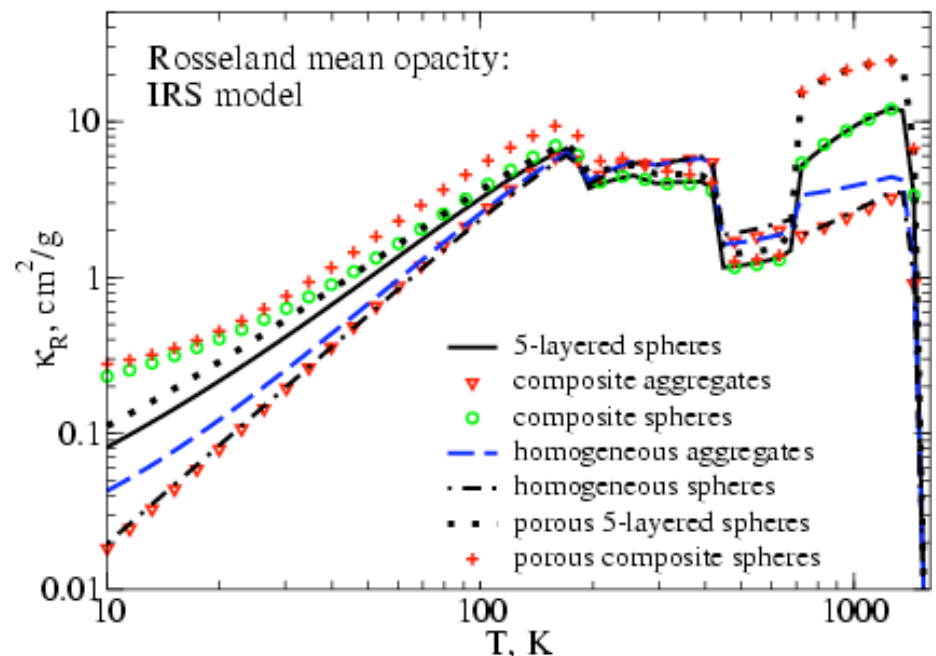
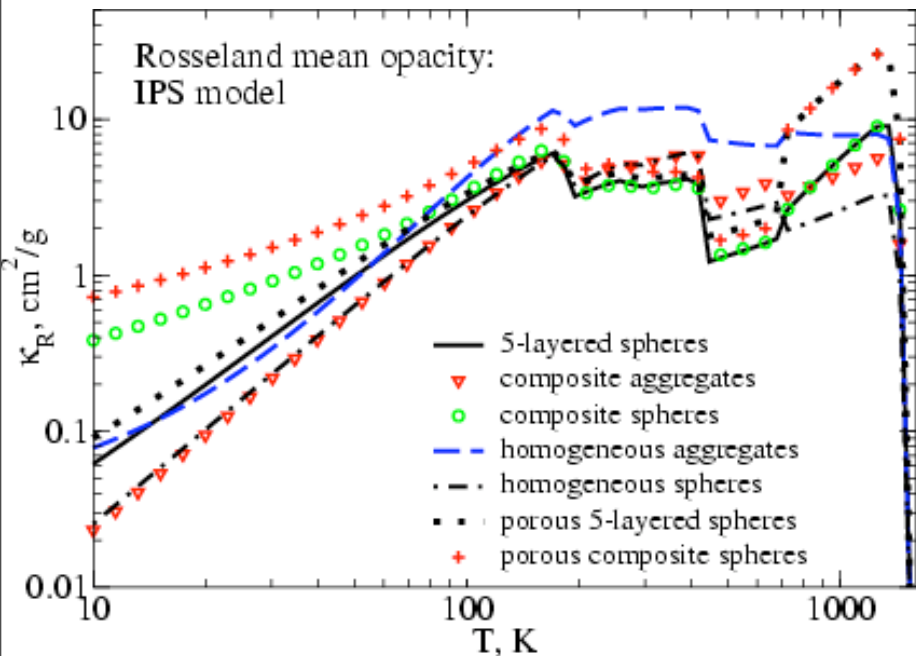
If  $10^5 \text{ K}$  is too hot, what supports the obscuring matter?

# Candidate Mechanisms

- Bouncing magnetized clouds stirred by orbital shear (K. & Begelman 1988, Beckert & Duschl 2004)  
But is this degree of elasticity plausible? And their collision rate must not be  $\gg$  the orbital frequency
- Warped thin disk (Sanders et al. 1989)  
But well-formed ionization cones are seen close to the center, and IR interferometry shows a thick structure at  $\sim 1$  pc in NGC 1068
- Magnetic wind (Königl & Kartje 1994)  
But origin of large-scale field? And mass-loss rate can be very large:  $10 N_{\text{H}24} (\text{h/r})(v_{r100})r_{\text{pc}} \text{ Msun/yr}$

# Another Candidate Mechanism: Radiation Pressure (Pier & K. 1992)

- If thermal continuum is created by dust reprocessing, there must be a large radiation flux through the obscuration
- $\kappa_{\text{midIR}}(\text{dust}) \sim 10\text{--}30 \kappa_T$  (Semenov et al. 2003)  
so  $(L/L_E)_{\text{eff}} \sim (10\text{--}30)L/L_E$



# Radiation Transfer and Dynamics Must Be Consistent

- $F_{\text{rad}}$  moves dusty gas, altering radiation transfer.
- New transfer solution changes  $F_{\text{rad}}$

# The Simplest Self-Consistent Picture

## Assumptions:

- Hydrostatic
- 2-d, axisymmetric, time-steady
- opacity independent of  $T$ , diffusion approximation (smooth density distribution)
- no sources of infrared inside the torus
- $1/l_{\text{Kep}} = j(r)$

## Solution: Entirely Analytic

Step 1: Hydrostatic balance + diffusion equilibrium lead to

$$r \, dj^2/dr + 2(1 - \alpha) j^2 = 3 - 2\alpha \quad \text{for } \alpha = -d \ln \Omega / d \ln r,$$

so that  $j^2(r) = [j_{in}^2 + f(\alpha)](r/r_{in})^{2(\alpha-1)} - f(\alpha)$

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In other words,

If  $F_{\text{rad},z} \sim F_{\text{grav},z}$  in a geometrically thick disk, then

$F_{\text{rad},r} \sim F_{\text{grav},r}$  likewise

So the orbiting matter must have sub-Keplerian rotation if it is to remain in equilibrium; magnetic angular momentum redistribution?



Step 2: Requiring both components of force balance to give a consistent density leads to

$$(\partial E / \partial z) / (z \Omega^2) - (\partial E / \partial r) / \{r \Omega^2 [1 - j^2(r)]\} = 0$$

which is analytically solvable by characteristics:

$E = \text{constant}$  on (almost) elliptical surfaces

Step 3: Two boundary conditions:

A) To fix the quantity of available matter---

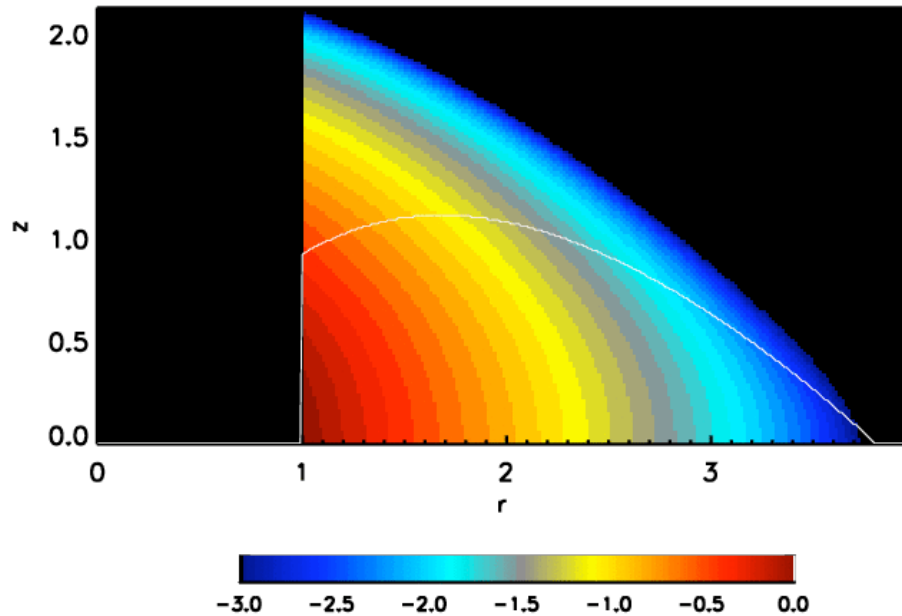
$$\rho(r, z=0) = (\tau_* / \kappa r_{\text{in}}) (r / r_{\text{in}})^{-\gamma}$$

B) To match the diffusion solution to its outgoing flux---

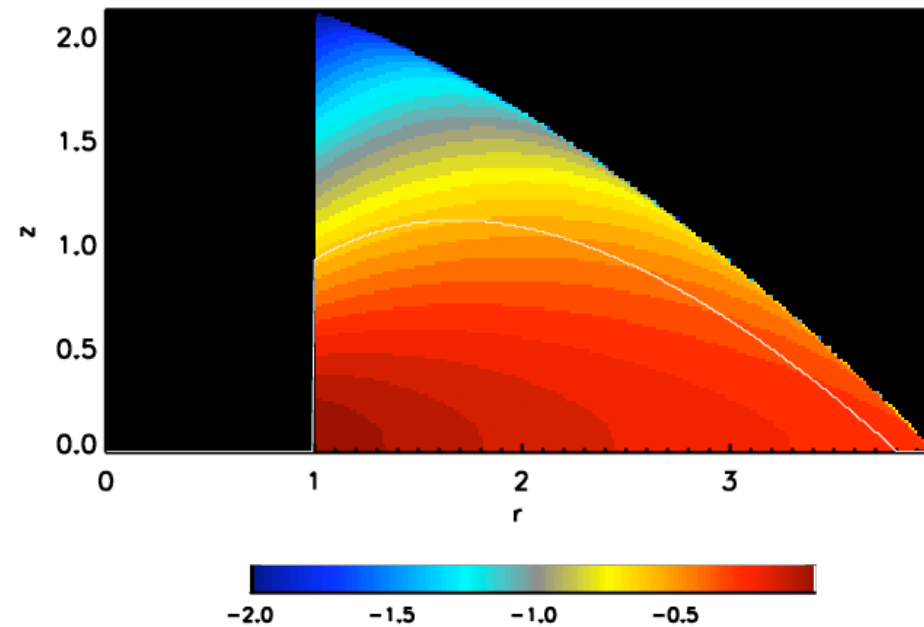
$F \sim cE$  at the dust photosphere (which determines  $\gamma$ )

# Example Solution

for  $\alpha = 3/2$ ,  $\tau_* = 10$ ,  $L/L_E = 0.1\text{---}0.3$

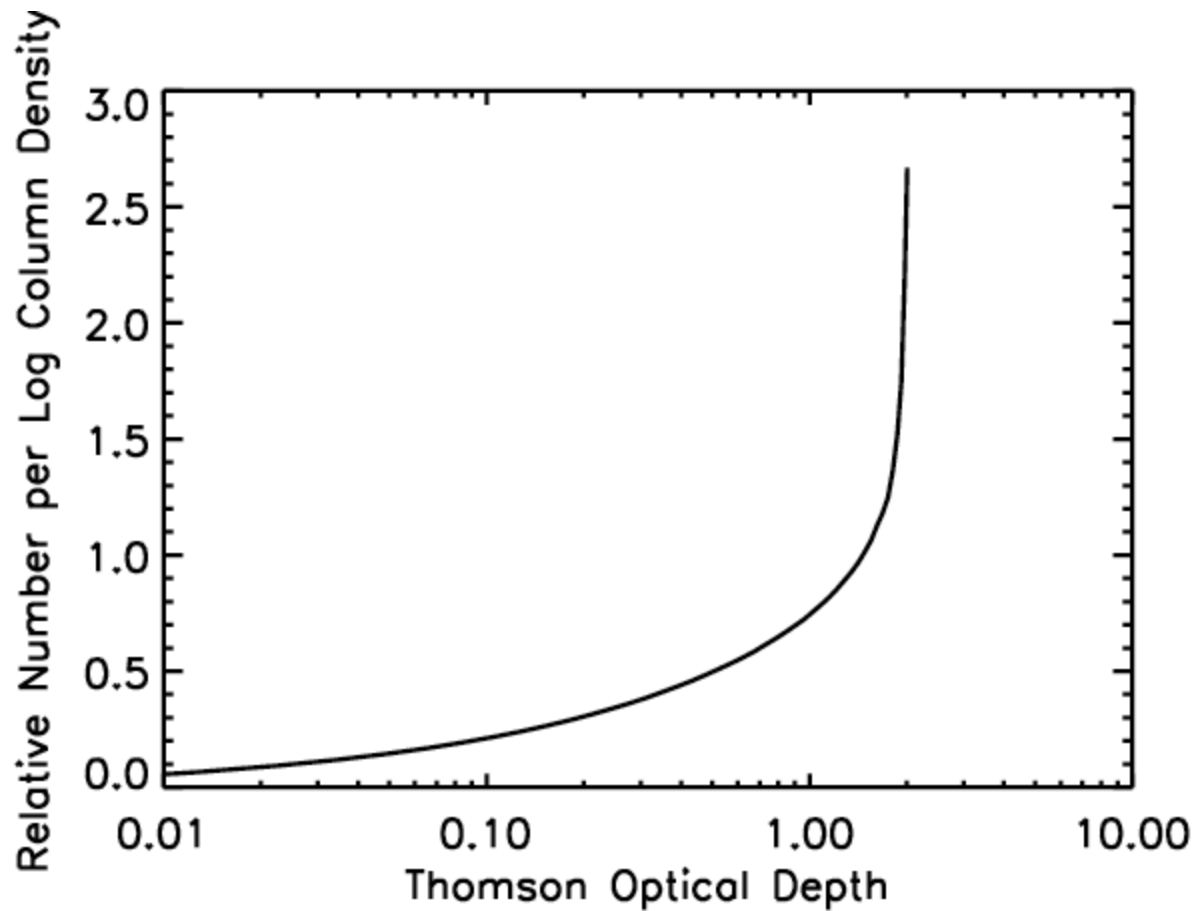


radiation energy density



gas density

# The X-ray Column Density Distribution Predicted by the Example Solution



## Conclusions

- Tori convert optical/UV flux to IR; there **must** therefore be a large IR flux through them
- Mid-IR opacity/mass  $\sim 10\text{—}30$  Thomson, **increasing** the effective  $F/F_E$  by that factor
- With some simplifying assumptions, a self-consistent hydrostatic equilibrium and 2-d diffusive transfer solution can be found
- The torus becomes geometrically thick when  $L/L_E \sim 0.03\text{—}0.3$  and the midplane  $\tau_T \sim 1$