

LETTER TO THE EDITOR

# Imaging the innermost circumstellar environment of the red supergiant WOH G64 in the Large Magellanic Cloud <sup>★</sup>

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## ABSTRACT

**Context.** Significant mass loss in the red supergiant (RSG) phase has great influence on the evolution of massive stars and their final fate as supernovae.

**Aims.** We present near-infrared interferometric imaging of the circumstellar environment of the dust-enshrouded RSG WOH G64 in the Large Magellanic Cloud.

**Methods.** WOH G64 was observed with the GRAVITY instrument at ESO's Very Large Telescope Interferometer (VLTI) at 2.0–2.45  $\mu\text{m}$ . We succeeded in imaging the innermost circumstellar environment of WOH G64 – the first interferometric imaging of an RSG outside the Milky Way.

**Results.** The reconstructed image reveals elongated compact emission with a semimajor and semiminor axis of  $\sim 2$  and  $\sim 1.5$  mas ( $\sim 13$  and  $9 R_*$ ), respectively. The GRAVITY data show that the stellar flux contribution at 2.2  $\mu\text{m}$  at the time of our observations in 2020 is much lower than predicted by the optically and geometrically thick dust torus model based on the VLTI/MIDI data taken in 2005 and 2007. We found a significant change in the near-infrared spectrum of WOH G64: while the (spectro)photometric data taken at 1–2.5  $\mu\text{m}$  before 2003 show the spectrum of the central RSG with H<sub>2</sub>O absorption, the spectra and *JHK'* photometric data taken after 2016 are characterized by a monotonically rising continuum with very weak signatures of H<sub>2</sub>O. This spectral change likely took place between December 2009 and 2016. On the other hand, the mid-infrared spectrum obtained in 2022 with VLT/VISIR agrees well with the spectra obtained before 2007.

**Conclusions.** The compact emission imaged with GRAVITY and the near-infrared spectral change suggest the formation of hot new dust close to the star, which gives rise to the monotonically rising near-infrared continuum and the high obscuration of the central star. The elongation of the emission may be due to the presence of a bipolar outflow or effects of an unseen companion.

**Key words.** infrared: stars – techniques: interferometric – stars: imaging – (stars:) supergiants – (stars:) circumstellar matter – stars: individual: WOH G64

## 1. Introduction

Significant mass loss in the red supergiant (RSG) phase is of great importance for the evolution of massive stars before they end their life in a supernova (SN) explosion. The RSGs at advanced evolutionary stages experience drastic mass loss with a mass-loss rate as high as  $\sim 10^{-4} M_{\odot} \text{ yr}^{-1}$  (Goldman et al. 2017). Recent analyses of very early-phase spectra of SNe – taken within a day after the explosion – suggest significant increases in the mass-loss rate in the RSG phase before the SN explosion (e.g., Yaron et al. 2017; Moriya et al. 2018; Zhang et al. 2023).

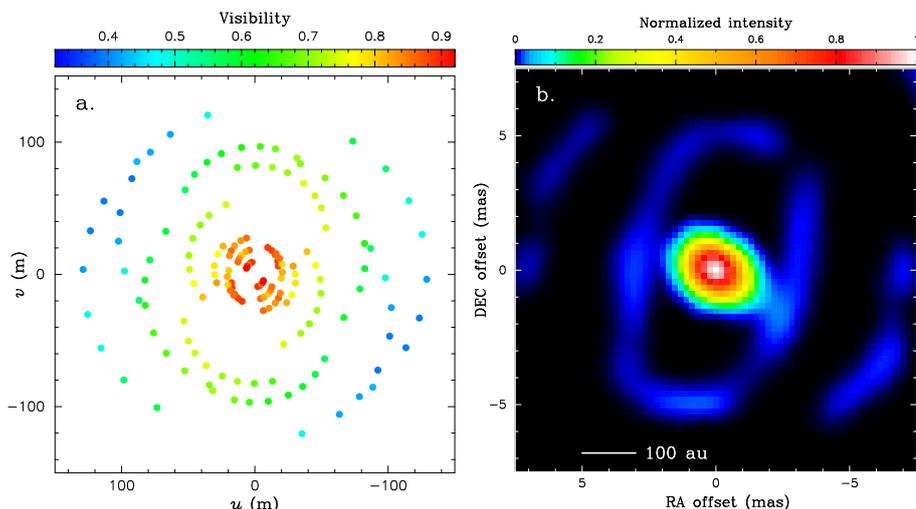
High-spatial-resolution observations of some RSGs reveal salient deviations from spherical symmetry in their circumstellar environment (e.g., Wittkowski et al. 1998; Monnier et al. 2004; Humphreys et al. 2007; O’Gorman et al. 2015). Nonspherical

mass loss can also be exemplarily seen in the dust ring around SN1987A, which is considered to have been shed in the RSG phase before the progenitor evolved into a blue supergiant and exploded (Crotts & Heathcoat 1991). Given the high multiplicity rate among massive stars (Mason et al. 2009; Sana et al. 2012), the asymmetric, enhanced mass loss in the RSG phase, which can be driven by binary interaction (e.g., Ercolino et al. 2024; Landri & Pejcha 2024), is essential not only for better understanding the evolution of massive stars but also for interpreting early-phase SN spectra.

The RSGs in the Large Magellanic Cloud (LMC) have the great advantage that their distances are much better known (50 kpc, Pietrzyński et al. 2013) compared to those of their Galactic counterparts. WOH G64 is the brightest RSG in the mid-infrared in the LMC, exhibiting a huge infrared excess with a high mass-loss rate on the order of  $10^{-4} M_{\odot} \text{ yr}^{-1}$  (Goldman et al. 2017). For this reason, it has been a subject of multiwavelength studies from the visible to the radio (e.g., van Loon et al. 1996; Levesque et al. 2009; Matsuura et al. 2016). Ohnaka et al. (2008) succeeded in spatially resolving the circumstellar dust environment of WOH

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**Fig. 1.** Visibility and image of WOH G64 obtained from our VLTI/GRAVITY observations. **a:**  $uv$  coverage of our VLTI/GRAVITY observations of WOH G64 with the calibrated visibility color-coded. North is up, east is to the left. **b:** Image of WOH G64 reconstructed at  $2.2 \mu\text{m}$  (with a spectral window of  $0.2 \mu\text{m}$ ) using IRBis with the maximum entropy regularization. North is up, east is to the left.

35 G64 using the mid-infrared interferometric instrument MIDI at  
 36 the Very Large Telescope Interferometer (VLTI). Their 2D radiative transfer modeling shows that the observed spectral energy distribution and the visibilities measured at  $8\text{--}13 \mu\text{m}$  can be simultaneously reproduced by a geometrically and optically thick torus viewed nearly pole-on. Moreover, the luminosity of  $\sim 2.8 \times 10^5 L_{\odot}$  derived from the dust torus model brought WOH G64 to fair agreement with the evolutionary track with an initial mass of  $25 M_{\odot}$ .

44 However, because MIDI was a two-telescope interferometer, it was not possible to obtain an image of WOH G64. In this Letter, we present the first infrared interferometric imaging of WOH G64.

## 48 2. Observations and data reduction

49 We carried out interferometric observations of WOH G64 with GRAVITY (GRAVITY Collaboration 2017) at VLTI on December 15 and 26, 2020 (UTC), at  $2.0\text{--}2.45 \mu\text{m}$ , using the Auxiliary Telescope (AT) configurations A0-G1-J2-J3 and A0-B2-C1-D0 with a maximum projected baseline length of 129 m (Program ID: 106.219D.001/002, P.I.: K. Ohnaka). We also observed HD32956 (F0/2IV/V, uniform-disk diameter =  $0.17 \text{ mas}$ , JMMC catalog: Bourges et al. 2017) and HD37379 (F6/7V, uniform-disk diameter =  $0.16 \text{ mas}$ ) for the interferometric and spectroscopic calibration. Our GRAVITY observations are summarized in Table A.1.

60 To increase the signal-to-noise (S/N) of the results, the raw GRAVITY data taken with a spectral resolution of 500 were first spectrally binned using a running box car filter, which resulted in a spectral resolution of 330. The spectral binning was applied to both the science target and the calibrators as well as the raw calibration files needed to create the P2VM. The spectrally binned raw data were then reduced with the GRAVITY pipeline ver 1.4.1<sup>1</sup>. The errors in the calibrated visibilities of WOH G64 were computed from the errors given by the pipeline and the variations in the transfer function calculated from all the calibrators observed on each night.

71 Figure 1a shows the  $uv$  coverage of our GRAVITY observations of WOH G64 with the visibility color-coded. The visibility falls off more rapidly in the northeast-southwest direction than in the northwest-southeast direction, which suggests that the object appears larger in the northeast-southwest direction. As Figs. B.1

and B.2 show, the visibilities and closure phases show no trace of the CO bands, although the  $2.3 \mu\text{m}$  band head is weakly seen in the spectrum (Fig. 3b). We reconstructed images from the GRAVITY data using IRBis (Hofmann et al. 2014) and MiRA (Thiébaud 2008). IRBis selects the best reconstruction based on the fit to the measurements and the distribution of the positive and negative residuals between the measured quantities (visibilities and closure phases) and those of the reconstructed image. The best reconstruction with IRBis was obtained with the maximum entropy regularization using the cost functions 1 and 2 defined in Hofmann et al. (2014). For the reconstruction with MiRA, two different regularizations (pixel difference quadratic and pixel intensity quadratic) were employed, and the optimal value of the hyperparameter was determined with the L-curve method. The spectrum of WOH G64 extracted from the GRAVITY data was spectroscopically calibrated, as is described in Appendix C.

We also obtained  $J$ -,  $H$ -, and  $K'$ -band photometric data on August 11, 2024, with the REMIR camera at the Rapid Eye Mount (REM) Telescope<sup>2</sup> (Molinari 2019). Details of the REM observations are described in Appendix D.

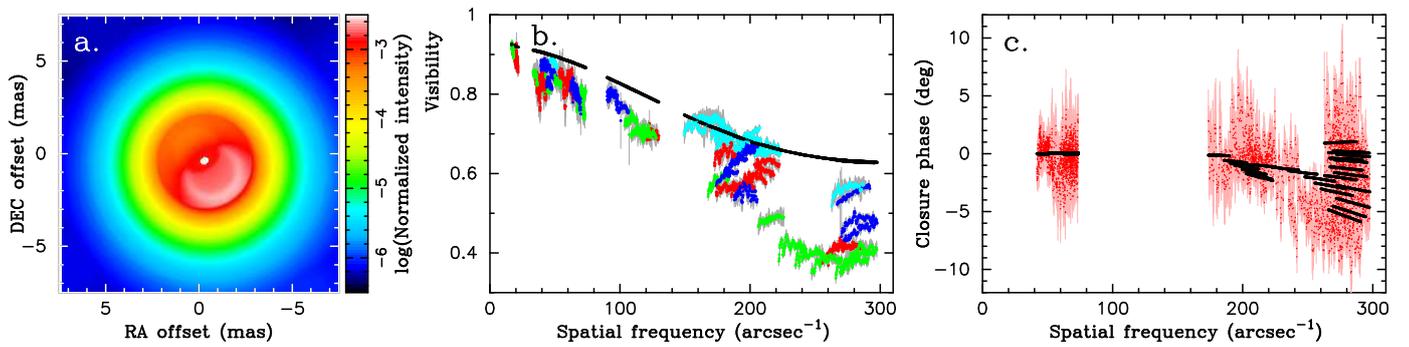
## 97 3. Results

98 Figure 1b shows the image reconstructed at  $2.2 \mu\text{m}$  with IRBis using the data between  $2.1$  and  $2.3 \mu\text{m}$ . The resolution of the image is  $\sim 1 \text{ mas}$ . Figure E.1 shows comparisons of the measured visibilities and closure phases with those computed from the IRBis reconstructed image. The image reconstructed with MiRA is shown in Fig. E.2, in which only the image obtained with the pixel difference quadratic regularization is presented, because the result obtained with the pixel intensity quadratic regularization is similar.

107 Both images obtained with IRBis and MiRA reveal compact elongated emission with a major and minor axis of  $\sim 4 \text{ mas}$  and  $3 \text{ mas}$ , respectively, with a position angle of the major axis of  $\sim 56^\circ$ . This corresponds to the position angle dependence of the visibility seen in Fig. 1a. The stellar radius of  $1730 R_{\odot}$  derived by Ohnaka et al. (2008) and the distance of  $50 \text{ kpc}$  translate into a stellar angular radius of  $0.16 \text{ mas}$ . Therefore, the semimajor and semiminor axes of the elongated emission correspond to  $\sim 13$  and  $\sim 9 R_{\star}$ , respectively. The central star does not clearly appear as a point source in our reconstructed image. This could be because

<sup>1</sup> <https://www.eso.org/sci/software/pipelines/gravity/>

<sup>2</sup> <http://www.rem.inaf.it>



**Fig. 2.** Comparison of the dust torus model of Ohnaka et al. (2008) with the GRAVITY data. **a:** Image predicted by the dust torus model at  $2.2 \mu\text{m}$  with an inclination angle of  $20^\circ$ . Note that the central star’s intensity (white dot at the center) is at least  $\sim 400$  times higher than that of the torus. North is up, east is to the left. **b:** Comparison of the visibility between the dust torus model and the GRAVITY data. The black dots (they appear to be solid lines due to the high density) represent the  $2.2 \mu\text{m}$  visibilities predicted by the dust torus model. The dots with the error bars represent the GRAVITY data with the position angle (PA) of the projected baseline color-coded: red ( $0^\circ \leq \text{PA} < 45^\circ$ ), green ( $45^\circ \leq \text{PA} < 90^\circ$ ), blue ( $90^\circ \leq \text{PA} < 135^\circ$ ), and light blue ( $135^\circ \leq \text{PA} < 180^\circ$ ). **c:** Comparison of the closure phase between the dust torus model and the GRAVITY data. The model closure phases are shown in black, while the measurements are plotted with the red dots with the error bars.

117 the central star is very faint or because it is entirely obscured.  
 118 The visibilities measured at the longest baselines tend to be flat  
 119 as a function of spatial frequency (Figs. 2b and E.1a), which can  
 120 be explained by the presence of a point source. However, data at  
 121 even longer baselines are needed to confirm this and derive the  
 122 fractional flux contribution of the central star in the  $K$  band.

123 A faint elliptical ring-like structure ( $\lesssim 3\%$  of the peak intensity)  
 124 with a semimajor and semiminor axis of  $\sim 5$  mas and  $3$  mas  
 125 ( $\sim 31$  and  $19 R_\star$ ), respectively, can be seen in dark blue in both  
 126 images reconstructed with IRBis and MiRA. This ring-like structure  
 127 might be interpreted as the inner rim of a dust disk or torus  
 128 viewed from an intermediate inclination angle, higher than in the  
 129 torus model of Ohnaka et al. (2008). The ring’s appearance in the  
 130 images reconstructed with different algorithms (MiRA and IR-  
 131 Bis) lends support to it being real. However, as Fig. E.2 shows,  
 132 the location of the ring is just inside the side lobe of the dirty  
 133 beam, and therefore we cannot entirely exclude the possibility  
 134 that it might be an artifact of the image reconstruction inherent  
 135 in the imperfect  $uv$  coverage. We refrain from further discussing  
 136 this structure until we can confirm it with better  $uv$  coverage.

137 Figure 2a shows the image predicted at  $2.2 \mu\text{m}$  from the dust  
 138 torus model of Ohnaka et al. (2008). The model image is charac-  
 139 terized by the central star and the emission from the inner rim on  
 140 the far side (light red to white region in the southwest). To exam-  
 141 ine whether the model is consistent with the GRAVITY data, we  
 142 show comparisons of the visibility and closure phase in Figs. 2b  
 143 and 2c. While the model closure phase is in fair agreement with  
 144 the measurements, the model predicts the  $2.2 \mu\text{m}$  visibility to  
 145 be too high compared to the data. We computed  $2.2 \mu\text{m}$  images  
 146 at different inclination angles within its uncertainty ( $0$ – $30^\circ$ )  
 147 and different position angles in the plane of the sky, but they all show  
 148 too high visibilities and too little elongation compared to the  
 149 data. This is because the model predicts the stellar flux contribu-  
 150 tion at  $2.2 \mu\text{m}$  to be too high compared to the GRAVITY data  
 151 (note that the central star’s intensity is higher than the intensity  
 152 of the light red to white region in the southwest by a factor of at  
 153 least  $\sim 400$ ).

#### 154 4. Change in the spectral shape in the infrared

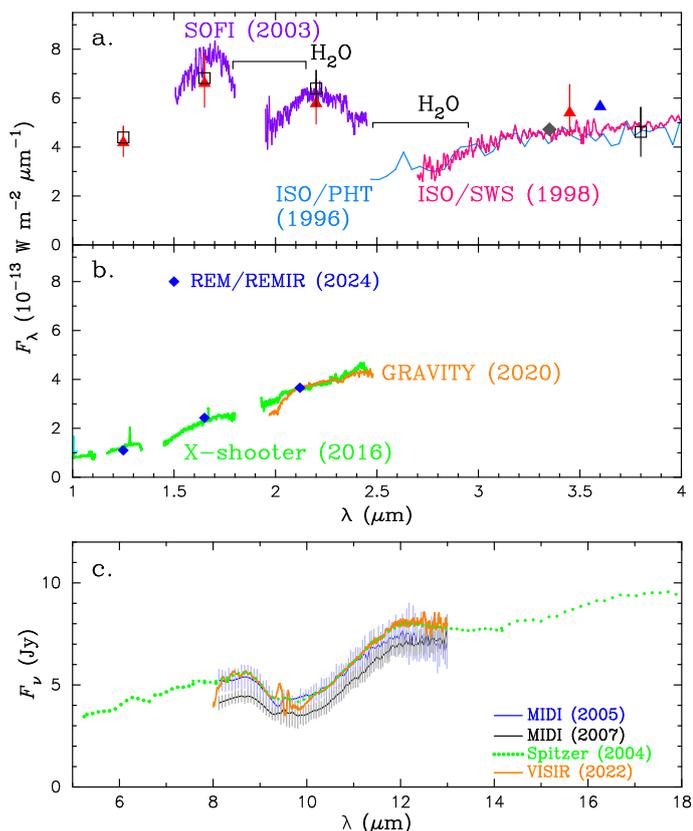
155 We found evidence that the difference in the  $2.2 \mu\text{m}$  stellar flux  
 156 contribution could be due to a systematic change in the circum-  
 157 stellar environment of WOH G64 between the MIDI observa-

158 tions in 2005 and 2007 and our GRAVITY observations in 2020.  
 159 Figure 3a shows (spectro)photometric data of WOH G64 taken  
 160 before 2010 collected from the literature (Wood et al. 1992;  
 161 Whitelock et al. 2003; Srinivasan et al. 2009, *Spitzer* Space Tele-  
 162 scope; Cutri et al. 2014, Wide-field Infrared Survey Explorer  
 163 (WISE); Vandebussche et al. 2002, Infrared Space Observatory  
 164 (ISO)/Short Wavelength Spectrometer (SWS); Trams et al. 1999,  
 165 ISO/Infrared Photo-polarimeter (PHT)). We also derived an  $HK$ -  
 166 band spectrum from the archival SOFI data taken with a spectral  
 167 resolution of 1000 (the reduction of the SOFI data is described  
 168 in Appendix F). The spectrum obtained at  $3.95$ – $4.10 \mu\text{m}$  with  
 169 ISAAC in 2001 (Matsuura et al. 2005) is flat, without noticeable  
 170 signatures of the  $\text{SiO}$  bands<sup>3</sup>, consistent with the ISO/SWS data.  
 171 These (spectro)photometric data show the spectrum of the RSG  
 172 with  $\text{H}_2\text{O}$  absorption bands, as is indicated in Fig. 3a.

173 On the other hand, Fig. 3b shows near-infrared spectra  
 174 of WOH G64 taken more recently with VLT/X-shooter (Pro-  
 175 gram ID: 097.D-0605(A), P.I.: S. Goldman; see Appendix G  
 176 for details of the X-shooter data and GRAVITY as well as the  
 177 REM/REMIR  $JHK'$  photometric data. Because the absolute flux  
 178 calibration of the X-shooter and GRAVITY spectra is uncertain,  
 179 we tentatively scaled them to the REM/REMIR  $K'$ -band flux just  
 180 to compare the spectral shape (the absolute flux scale likely var-  
 181 ied among the X-shooter, GRAVITY, and REM observations).  
 182 The X-shooter, GRAVITY, and REM/REMIR data all indicate  
 183 a monotonically rising continuum. This is in marked contrast to  
 184 the RSG spectrum seen in Fig. 3a. The  $\text{H}_2\text{O}$  absorption is also  
 185 much less pronounced in the X-shooter spectrum compared to  
 186 the SOFI data. The molecular spectral features in cool variable  
 187 stars like WOH G64 change with the variability phase and cycle.  
 188 The bars of the  $JHK$  photometric data in Fig. 3a show the  
 189 variability amplitudes measured from March 1995 to April 1998  
 190 by Whitelock et al. (2003). The variability amplitudes measured  
 191 from January 1987 to November 1991 by Wood et al. (1992) are  
 192 similar<sup>4</sup>. The rising continuum spectrum is difficult to explain  
 193 by the variability seen in the  $J$ ,  $H$ , and  $K$  bands.

<sup>3</sup> The absolute flux scale is uncertain by a factor of two, which is why the ISAAC spectrum was not included in Fig. 3

<sup>4</sup> The variabilities in the  $J$  and  $H$  bands measured by Wood et al. (1992) are not shown in Fig. 3a, because they overlap with those measured by Whitelock et al. (2003). Only the  $K$ - and  $L$ -band variabilities are indicated with the bars.



**Fig. 3.** Near- and mid-infrared spectrum of WOH G64. **a:** Spectrophotometric and photometric data taken at 1–4  $\mu\text{m}$  before 2010. The light blue, pink, and violet lines represent the spectra taken with ISO/PHT in 1996, ISO/SWS in 1998, and SOFI in 2003, respectively. The red triangles and open squares correspond to the photometric data from Whitelock et al. (2003) and Wood et al. (1992), respectively. The bars show the (peak-to-peak) variability amplitudes, not the measurement errors. The gray diamond and blue triangle represent the photometric data taken with WISE in 2009–2010 and *Spitzer* in 2005–2006, respectively. The  $\text{H}_2\text{O}$  absorption features in the RSG spectrum are also indicated. **b:** Spectra of WOH G64 taken more recently. The green and orange lines represent the spectra taken with X-shooter in 2016 and GRAVITY in 2020, respectively. The blue diamonds show the *JHK'* photometric data taken in 2024 with REM/REMIR. The GRAVITY and X-shooter data are scaled to the REM/REMIR *K'*-band flux. The spikes in the X-shooter spectrum as well as the very tiny peak at  $\sim 2.16 \mu\text{m}$  in the GRAVITY spectrum are the residual of the correction for the absorption lines in the spectrum of the calibrator. **c:** Mid-infrared spectra of WOH G64 showing the  $10 \mu\text{m}$  silicate feature in absorption. The blue and black lines represent the spectra obtained with MIDI in 2005 and 2007, respectively, by Ohnaka et al. (2008). The dotted green line shows the *Spitzer*/IRS spectrum taken in 2004, while the orange line corresponds to the VISIR spectrum taken in 2022. The wavelength range between 9.3 and  $9.8 \mu\text{m}$  is severely affected by the telluric absorption.

spectra within the errors. A comparison of the MIDI and *Spitzer* spectra with other data in the literature presented in Ohnaka et al. (2008, their Fig. 3) indicates that the long-term variation in the mid-infrared spectrum of WOH G64 is small.

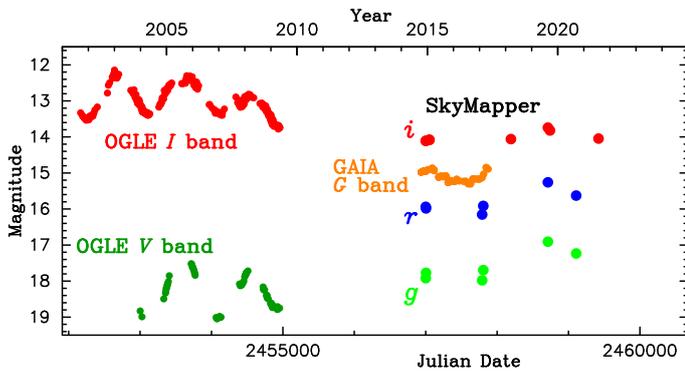
## 5. Discussion and conclusion

The results that the noticeable change in the spectral shape is seen in the near-infrared but not in the mid-infrared may be explained by hot dust forming close to the star. Haubois et al. (2019, 2023) detected dust clouds forming at  $\sim 1.5 R_\star$  around the optically bright RSG Betelgeuse (which is much less dust-enshrouded than WOH G64) by near-infrared and visible polarimetry. They concluded that transparent grains such as  $\text{Al}_2\text{O}_3$  and Fe-free silicates ( $\text{MgSiO}_3$  and  $\text{Mg}_2\text{SiO}_4$ ) can explain the observed data, although they favor the Fe-free silicates over  $\text{Al}_2\text{O}_3$ . These grains condense close to the star at  $\sim 1500 \text{ K}$ . Dynamical modeling of winds of asymptotic giant branch stars shows that Fe-rich silicates, whose absorption efficiency is much higher than the aforementioned species, condense onto the transparent grains a little farther out (Höfner et al. 2022). If this also occurs in RSGs like WOH G64, and the grains with mantles of Fe-rich silicates are optically thick in the near-infrared, it could account for the monotonically rising continuum spectrum seen in the near-infrared and the high obscuration of the central star.

The absence of a noticeable change in the mid-infrared spectrum suggests that the hot dust should be confined in a region close to the star. This is in qualitative agreement with the semi-major and semiminor axes of the elongated emission of  $\sim 13$  and  $\sim 9 R_\star$ , respectively. The elongated emission may be due to a bipolar outflow along the axis of the dust torus. Goldman et al. (2017) interpreted the 1612 MHz OH maser emission of WOH G64 at an expansion velocity of  $23.8 \text{ km s}^{-1}$  as spherical expansion or a bipolar outflow along our line of sight. Alternatively, the elongation may be caused by the interaction with an unseen companion. The torus model of Ohnaka et al. (2008) is characterized by an inner radius of  $15 R_\star$  ( $2.4 \text{ mas}$ ), in which the elongated emission with  $13 \times 9 R_\star$  ( $2 \times 1.5 \text{ mas}$ ) can fit, although this cannot be taken as definitive evidence for a companion. Levesque et al. (2009) discussed the possibility of an unseen hot companion and proposed spectroscopy in the blue to prove or disprove the hypothesis. However, neither positive nor negative detections have been reported in the literature.

While the SOFI data taken in 2003 clearly shows the  $\text{H}_2\text{O}$  absorption in the RSG spectrum, the X-shooter spectrum taken in 2016 shows the rising continuum. Therefore, the spectral shape change in the near-infrared should have started at some epoch between 2003 and 2016. Although no near-infrared spectra taken between 2003 and 2016 can be found in the literature, we can set an additional constraint on when the spectral change started as follows. Levesque et al. (2009) derived an extinction toward the central star  $A_V = 6.82 \approx \tau_V$  by fitting the visible spectrum of WOH G64 taken in December 2009. This optical depth in the visible translates into  $\tau_K = 0.62$  if we adopt the complex refractive index of the warm silicate of Ossenkopf et al. (1992) with the grain size ( $a$ ) distribution,  $\propto a^{-3.5}$ , between  $0.005$  and  $0.1 \mu\text{m}$ , as Ohnaka et al. (2008) assumed ( $\tau_K$  is lower,  $0.18$ , for a single grain size of  $0.1 \mu\text{m}$ ). With  $\tau_K < 1$ , the central star should still have been visible in the *K* band in December 2009. Therefore, it is likely that the formation of hot dust started at some epoch between December 2009 and 2016, and had not yet taken place at the time of our MIDI observations in 2005 and 2007.

The formation of new hot dust also means that the central star is now more obscured than the epochs before 2009. We collected



**Fig. 4.** Visible light curves of WOH G64 from 2001 to 2021. The red and dark green dots represent the photometric measurements in the  $I$  and  $V$  bands, respectively, from the OGLE project. The orange dots correspond to the *Gaia* data in the  $G$  band. The red, blue, and light green dots represent the photometric data from the SkyMapper Southern Survey in the  $i$ ,  $r$ , and  $g$  bands, respectively.

266 visible photometric data from the OGLE project (Soszyński et al. 2009) and SkyMapper Southern Survey (Onken et al. 2024) and show the visible light curves of WOH G64 in Fig. 4. The  $V$ - and  $I$ -band light curves obtained by the OGLE project show the periodic variation until the middle of 2009. There are noticeable differences in the filter systems between OGLE and SkyMapper. However, the OGLE  $I$  band and SkyMapper  $i$  band are relatively close, with central wavelengths of 801 nm (OGLE  $I$  band) and 779 nm (SkyMapper  $i$  band). The SkyMapper  $i$ -band data suggest that WOH G64 is fainter after the end of 2014 than in the years covered by the OGLE data. Still, the current sparsity of the data in the literature – particularly between 2010 and 2015, exactly when the formation of hot dust likely started – makes it difficult to examine whether WOH G64 appeared much fainter at the time of our GRAVITY observations in 2020 than in the past.

281 It is also worth noting that WOH G64 seems to be brighter in the  $g$  band than in the  $V$  band. Given that the  $g$  band is centered at 510 nm, shorter than the OGLE  $V$  band at 540 nm, and the visible flux of WOH G64 steeply decreases at shorter wavelengths (van Loon et al. 2005), we do not expect the  $g$ -band flux to be higher than the  $V$ -band flux. This high  $g$ -band flux may be due to the scattering by the transparent grains forming close to the star, as was proposed above, although this is not conclusive. Continuous long-term multiband photometric monitoring covering at least the pulsation period of 840–930 days as well as visible spectroscopic monitoring is necessary to confirm or refute the effects of the new dust formation in the visible.

293 While the mid-infrared spectrum of WOH G64 shows little long-term change, it is interesting to probe whether there has been a change in the structure of the circumstellar environment on spatial scales larger than the new dust formation region close to the star. The VLTI/MATISSE instrument (Lopez et al. 2022) allows us to obtain spectro-interferometric data in the  $L$  (3–4.2  $\mu\text{m}$ ),  $M$  (4.5–5.0  $\mu\text{m}$ ), and  $N$  bands (8–13  $\mu\text{m}$ ). MATISSE observations and radiative transfer modeling will set much tighter constraints on its circumstellar environment than MIDI did.

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396 **Appendix A: Observation log**

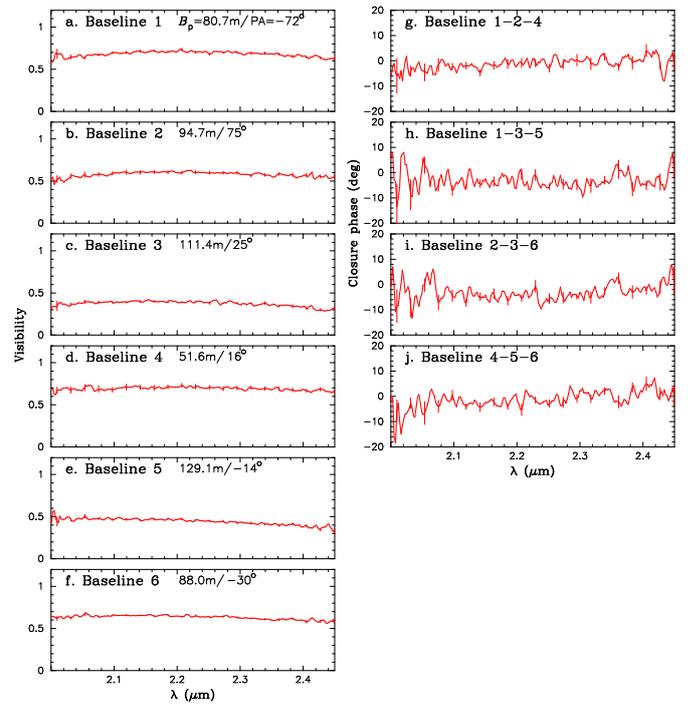
 397 Table A.1 shows the summary of our GRAVITY observations of  
 398 WOH G64.

**Table A.1.** Our VLTI/GRAVITY observations of WOH G64.

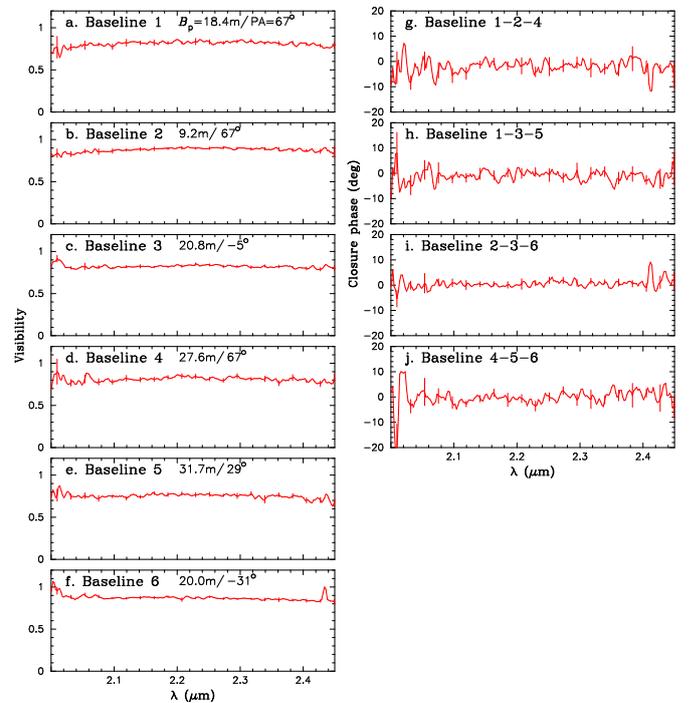
#	$t_{\text{obs}}$ UTC	$B_p$ (m)	PA ( $^\circ$ )	Seeing ( $''$ )	$\tau_0$ (ms)	$N_{\text{exp}}$
WOH G64						
DIT = 30 s, $N_f = 12$						
2020 December 15						
AT configuration: A0-G1-J2-J3						
1	01:57:23	70.9/96.1/123.4/ 56.3/126.1/79.6	-45/-81/59/ 52/26/8	0.57	6.2	1
2	02:42:02	74.3/96.8/121.3/ 55.5/127.7/82.4	-52/-87/50/ 43/16/-2	0.43	6.6	2
3	03:38:15	77.7/96.6/117.7/ 54.1/128.9/85.4	-61/85/39/ 31/3/-15	0.55	6.5	2
4	04:46:15	80.7/94.7/111.4/ 51.6/129.1/88.0	-72/75/25/ 16/-14/-30	0.38	6.6	1
5	05:39:33	81.9/92.0/105.4/ 49.2/128.2/89.2	-81/67/14/ 4/-26/-42	0.31	6.5	1
6	06:30:36	82.4/88.4/98.4/ 46.4/126.7/89.9	-89/60/2/ -8/-38/-53	0.32	7.5	2
7	07:30:54	82.0/83.0/89.8/ 43.0/124.8/90.3	82/51/-11/ -23/-53/-67	0.34	6.9	2
2020 December 26						
AT configuration: A0-B2-C1-D0						
8	01:13:48	18.8/9.4/19.9/ 28.2/32.0/18.6	77/77/8/ 77/42/-20	0.77	3.9	1
9	02:11:48	18.4/9.2/20.8/ 27.6/31.7/20.0	67/67/-5/ 67/29/-31	0.72	4.3	1
10	04:20:12	16.8/8.4/22.1/ 25.1/30.3/22.1	45/45/-34 45/-1/-56	0.70	4.8	1
11	05:04:03	15.9/7.9/22.3/ 23.8/29.4/22.5	37/37/-43/ 37/-12/-64	0.70	4.7	2

**Notes.**  $B_p$ : Projected baseline length. PA: Position angle of the baseline vector projected onto the sky. DIT: Detector Integration Time.  $N_f$ : Number of frames in each exposure.  $N_{\text{exp}}$ : Number of exposures. The seeing and the coherence time ( $\tau_0$ ) were measured in the visible.

 399 **Appendix B: Observed visibilities and closure**  
 400 **phases of WOH G64**

 401 Figures B.1 and B.2 show examples of the visibilities and closure  
 402 phases observed at the A0-G1-J2-J3 and A0-B2-C1-D0 config-  
 403 urations, respectively. The observed interferometric data (spec-  
 404 trally binned to the resolution of 330) do not show signatures of  
 405 the CO bands longward of 2.3  $\mu\text{m}$ .


**Fig. B.1.** Visibilities and closure phases of WOH G64 observed at the A0-G1-J2-J3 configuration. **a-f:** Visibility. The projected baselines ( $B_p$ ) and position angle (PA) are given in each panel. **g-j:** Closure phase. The baselines that form the telescope triplet are given in each panel.



**Fig. B.2.** Visibilities and closure phases of WOH G64 observed at the A0-B2-C1-D0 configuration, shown in the same manner as Fig. B.1.

## 406 Appendix C: Calibration of the GRAVITY spectra

407 We obtained the spectroscopically calibrated spectrum of WOH  
408 G64 as follows:

$$409 F_{\text{sci}}^{\text{true}} = \frac{\eta_{\text{cal}}}{\eta_{\text{sci}}} F_{\text{sci}}^{\text{obs}} \times \frac{F_{\text{cal}}^{\text{true}}}{F_{\text{cal}}^{\text{obs}}},$$

410 where  $F_{\text{sci}(\text{cal})}^{\text{true}}$  and  $F_{\text{sci}(\text{cal})}^{\text{obs}}$  denote the true and observed spectra of  
411 the science target (sci) or the calibrator (cal), respectively.  $\eta_{\text{sci}(\text{cal})}$   
412 represents the fraction of the flux injected into the fibers of the  
413 beam combiner for the science target or calibrator. The spectrum  
414 was derived from the data without the spectral binning, because  
415 the S/N of the spectral data was sufficiently high. We used the  
416 calibrator HD37379 (F6/7V) for the spectrophotometric calibration.  
417 To approximate the true spectrum of HD37379, we used  
418 the flux-calibrated spectrum of HD126660 (F7V) taken with the  
419 InfraRed Telescope Facility (IRTF Spectral Library<sup>5</sup>, Rayner et  
420 al. 2009) because its spectral type and luminosity class are very  
421 close to those of HD37379. The IRTF spectrum taken with a  
422 spectral resolution of  $\lambda/\Delta\lambda = 2000$  was convolved to match the  
423 spectral resolution of 500 of our GRAVITY (spectral) data and  
424 then scaled to match the  $K$  magnitude of HD37379 by multi-  
425 plying by  $f_{K,\text{HD37379}}/f_{K,\text{HD126660}}$ , where  $f_{K,\text{HD37379}(\text{or HD126660})}$   
426 denotes the  $K$ -band flux of HD37379 or HD126660. The spectrum  
427 obtained in this manner was used for  $F_{\text{cal}}^{\text{true}}$  in the above spectro-  
428 scopic calibration.

429 While the spectra obtained with four ATs agree well in the  
430 shape, the absolute flux calibration of the GRAVITY spectra is  
431 difficult. First, the fraction of the flux injected into the GRAV-  
432 ITY's beam combiner fibers depends on the performance of the  
433 adaptive optics (AO) system NAOMI of the ATs (Woillez et  
434 al. 2019). While the calibrator HD37379 is bright enough for  
435 NAOMI to work properly, the  $G$ -band magnitude of WOH G64  
436 is just at the limit of NAOMI ( $G = 15$ ), which resulted in de-  
437 graded AO performance for WOH G64 compared to HD37379.  
438 This means that the fraction of the flux injected into the fibers  
439 was systematically lower for WOH G64 than for HD37379. Fur-  
440 thermore, the fiber injection fraction also depends on other fac-  
441 tors such as the aberration and tip/tilt correction. Jovanovic et al.  
442 (2017) show that the coupling efficiency to single-mode fibers  
443 varies up to 30%. Given this large uncertainty, we tentatively  
444 scaled the GRAVITY spectra to the  $K'$ -band magnitude measured  
445 with REM/REMIR.

## 446 Appendix D: REM observations of WOH G64

447 The RSG WOH G64 was observed with the REM/REMIR camera  
448 with the  $J$ ,  $H$ , and  $K'$  filters on 2024 August 11 (UTC).  
449 We obtained 10 frames with each filter, using an exposure time  
450 of 15 s for each filter. After flat-fielding, sky subtraction, and  
451 averaging of the frames, the flux of WOH G64 was measured  
452 by PSF (point spread function) photometry and calibrated using  
453 five stars in the same field of view, for which the 2MASS  $JHK_s$   
454 magnitudes are available. The 2MASS  $K_s$  magnitudes were con-  
455 verted to the  $K'$  magnitudes using the relations given on the  
456 GEMINI web page<sup>6</sup> and the 2MASS web page<sup>7</sup>. The  $J$ -,  $H$ -,  
457 and  $K'$ -band fluxes of WOH G64 and their errors are listed in  
458 Table D.1.

**Table D.1.** Near-infrared flux of WOH G64 measured with  
REM/REMIR on 2024 August 11.

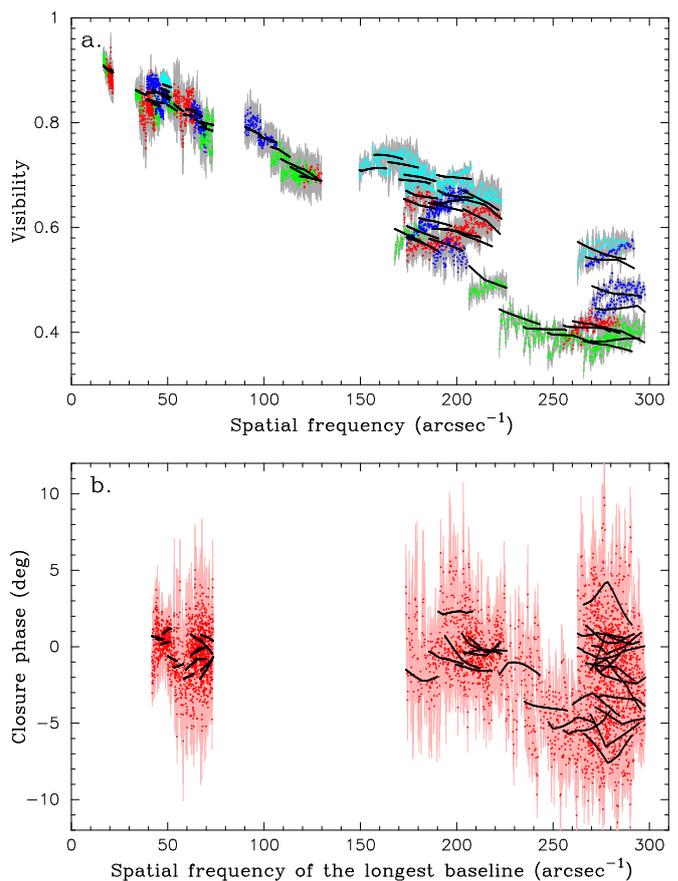
Band	Flux ( $10^{-13} \text{ W m}^{-2} \mu\text{m}^{-1}$ )
$J$	$1.10 \pm 0.04$
$H$	$2.43 \pm 0.07$
$K'$	$3.66 \pm 0.19$

## 459 Appendix E: Image reconstruction

460 Comparisons of the visibilities and closure phases computed  
461 from the IRBis reconstructed image and the GRAVITY data are  
462 shown in Fig. E.1.

463 The image reconstructed with MiRA using the pixel differ-  
464 ence quadratic regularization is shown in Fig. E.2, together with  
465 the image obtained with IRBis and the dirty beam. The MiRA  
466 image reconstructed with the pixel intensity quadratic regular-  
467 ization is very similar to the one obtained with the pixel dif-  
468 ference quadratic regularization. A flat prior was used in both  
469 reconstructions with MiRA.

470 We reconstructed images using the data with spectral win-  
471 dows centered at  $2.2 \mu\text{m}$  with different widths of 0.05, 0.1, and  
472  $0.2 \mu\text{m}$ . The reconstructed images appear to be very similar.  
473 Therefore, we show the image obtained with the width of  $0.2 \mu\text{m}$ ,  
474 which slightly enhances the  $uv$  coverage.

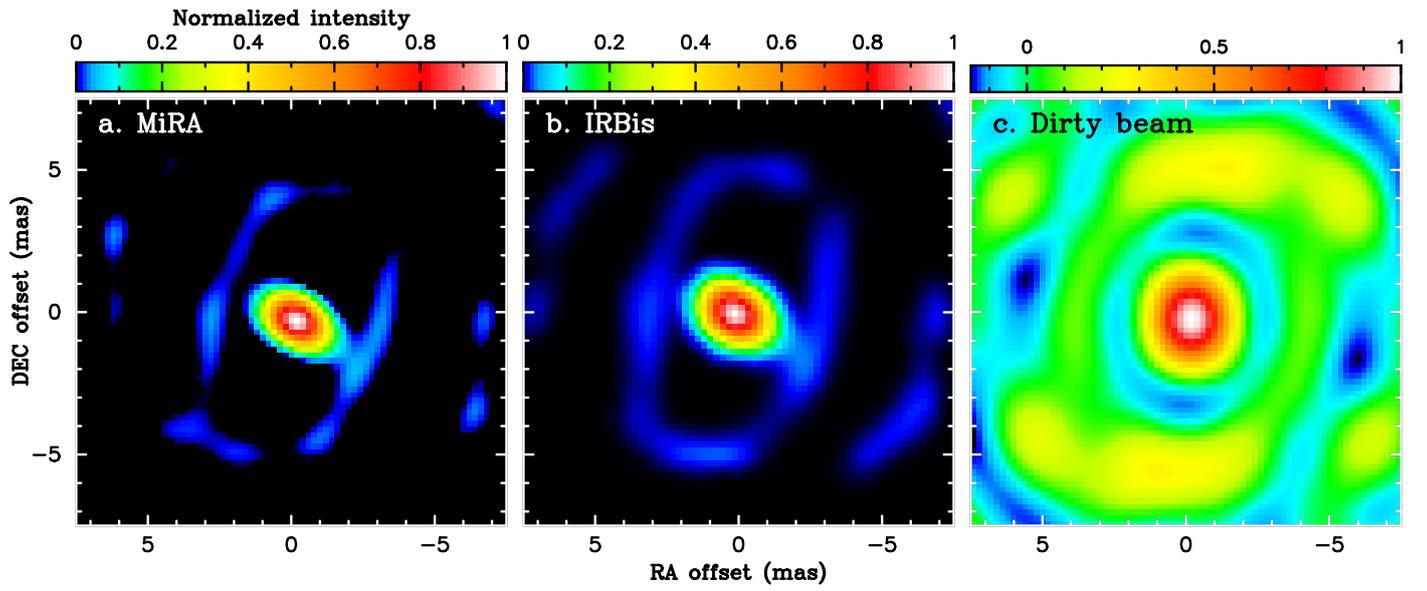


**Fig. E.1.** Comparison of the visibilities and closure phases of the recon-  
460 structed image shown in Fig. 1b with the GRAVITY measurements. **a:**  
461 Visibility. **b:** Closure phase. In either panel, the black dots (they appear  
462 to be solid lines due to the high density) represent the visibilities or clo-  
463 sure phases computed from the reconstructed image. The measurements  
464 are shown in the same manner as in Figs. 2b and 2c.

<sup>5</sup> [http://irtfweb.ifa.hawaii.edu/~spex/IRTF\\_Spectral\\_Library/](http://irtfweb.ifa.hawaii.edu/~spex/IRTF_Spectral_Library/)

<sup>6</sup> <https://www.gemini.edu/observing/resources/near-ir-resources/photometry/niri-filter-color-transformations>

<sup>7</sup> [https://irsa.ipac.caltech.edu/data/2MASS/docs/releases/allsky/doc/sec6\\_4b.html](https://irsa.ipac.caltech.edu/data/2MASS/docs/releases/allsky/doc/sec6_4b.html)



**Fig. E.2.** Images of WOH G64 reconstructed at  $2.2 \mu\text{m}$  with MiRA and IRBis. **a:** Image reconstructed using MiRA with the pixel difference quadratic regularization. **b:** Image reconstructed with IRBis (the same image as in Fig. 1b). **c:** Dirty beam. North is up, east to the left in all panels.

## 475 **Appendix F: SOFI observations of WOH G64**

476 We downloaded the SOFI archival data of WOH G64 and the  
 477 spectroscopic calibrator HIP21984 (G3V) taken on 2003 June  
 478 29, covering from 1.5 to 2.5  $\mu\text{m}$  with a spectral resolution of  
 479 1000 (Program ID: 71.B-0558(A), P.I.: J. Blommaert). The data  
 480 were processed with the SOFI pipeline ver. 1.5.12. The spectro-  
 481 scopic calibration of the data of WOH G64 was done in the same  
 482 manner as described in Appendix C but with both  $\eta_{\text{sci}}$  and  $\eta_{\text{cal}}$  set  
 483 to 1 because they are irrelevant for the SOFI observations. We  
 484 used HD10697 (G3Va) as a proxy star because of its similarity to  
 485 HIP21984 with respect to the spectral type and luminosity class.  
 486 The photometrically calibrated spectrum of HD10697 available  
 487 in the IRTF Spectral Library (Rayner et al. 2009) was convolved  
 488 to match the spectral resolution of 1000 of the SOFI data and  
 489 then multiplied by the  $K$ -band flux ratio  $f_{K,\text{HIP21984}}/f_{K,\text{HD10697}}$ .  
 490 This was used for  $F_{\text{cal}}^{\text{true}}$  in the spectrophotometric calibration.

## 491 **Appendix G: X-shooter observations of WOH G64**

492 We downloaded the reduced (phase 3) X-shooter spectra of  
 493 WOH G64 and the spectroscopic calibrator HIP28322 (B8V)  
 494 taken on 2016 July 27 (Program ID: 097.D-0605(A), P.I.: S.  
 495 Goldman). While the reduced data of WOH G64 are spectro-  
 496 scopically calibrated, the telluric lines are not removed. We at-  
 497 tempted to remove them as much as possible in the same man-  
 498 ner as described in Appendix C with both  $\eta_{\text{sci}}$  and  $\eta_{\text{cal}}$  set to 1.  
 499 To obtain  $F_{\text{cal}}^{\text{true}}$ , we used HD147550 (B9V) as a proxy star. Its  
 500 near-infrared spectrum available in the X-shooter Spectral Li-  
 501 brary (Gonneau et al. 2020), which is flux-calibrated and telluric-  
 502 corrected, was scaled to match the  $K$ -band flux of HIP28322  
 503 and used for  $F_{\text{cal}}^{\text{true}}$  in the spectroscopic calibration. The result-  
 504 ing spectrum was convolved to match the spectral resolution of  
 505 the GRAVITY spectrum of WOH G64. However, as mentioned  
 506 in Gonneau et al. (2020), the absolute flux calibration of the X-  
 507 shooter spectrum is difficult due to the slit loss. Therefore, we  
 508 tentatively scaled the X-shooter spectrum to match the  $K'$ -band  
 509 flux measured with REM/REMIR.

## 510 **Appendix H: VISIR observation of WOH G64**

511 Our VISIR observation of WOH G64 at 8–13  $\mu\text{m}$  took place  
 512 on 2022 October 7 (UTC). A slit width of 1'' was used, which  
 513 resulted in a spectral resolution of 300. The data of WOH G64  
 514 and the calibrator HD33554 were first reduced with the VISIR  
 515 pipeline ver. 4.4.2, and the spectrum of WOH G64 was calibrated  
 516 in the same manner as described in Appendix C with both  $\eta_{\text{sci}}$   
 517 and  $\eta_{\text{cal}}$  set to 1. We used the absolutely calibrated spectrum of  
 518 HD33554 presented in Cohen et al. (1999) as  $F_{\text{cal}}^{\text{true}}$ .