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# The formation of low-mass stars and brown dwarfs by disc fragmentation

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## The formation of low-mass stars and brown dwarfs

#### gravo-turbulent fragmentation of molecular clouds

Padoan & Nordland 2002; Bate et al. 2004; Goodwin et al. 2004 Hennebelle & Chabrier 2008, 2010

#### premature ejection of protostellar embryos

Clarke & Reipurth, Bate et al. 2003, 2004; Goodwin et al. 2004

#### disc fragmentation

Stamatellos, Hubber & Whitworth 2007, MNRAS Stamatellos & Whitworth 2009, MNRAS, Bate et al. 2003

## Brown dwarf and low-mass star formation by disc fragmentation

Stamatellos & Whitworth 2009, MNRAS, 392, 413 "The properties of brown dwarfs and low-mass hydrogen-burning stars formed by disc fragmentation"

✓ The shape the low-mass end of the IMF

✓ The presence of discs around brown dwarfs (BDs don't have to form like solar-mass stars in order to have discs!)

The formation of BD-BD binaries (produce both tight and wide binaries)
The formation of free-floating planetary-mass objects

✓ Specific low-mass binary characteristics: BDs that are companions to Sun-like stars are more likely (25-50%) to be in binaries than brown dwarfs in the field (~10%) (Burgasser et al. 2005; Faherty et al. 2010)

✓ The brown dwarf desert

- There are many planets and low-mass stars close (<5 AU) companions to Sun-like stars, but almost no brown dwarfs (Marcy & Butler 2000).
- The brown dwarf desert may extend out to ~1000 AU (Gizis et al. 2001) but is less "dry" of brown dwarfs outside ~50 AU (Neuhauser et al. 2003).



#### time ~ 5,000 yr



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#### time ~ 20,000 yr

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#### time ~ 200,000 yr



## Radiative feedback suppresses the formation of low-mass stars and brown dwarfs?

Offner et al. 2009; Bate 2009, 2011



■Almost no brown dwarfs form at all → something is missing

### The importance of radiative feedback from protostars

#### **Simulation I:**

Radia tive feedback is included down to the sink radius (1AU) [similar to Bate's 2009 simulations]

> Initial conditions: turbulent cloud core  $M = 5.4M_{\odot}$   $N_{\rm SPH} = 10^6$  particles

$$\rho(r) = \frac{\rho_{\text{kernel}}}{(1 + (r/R_{\text{kernel}})^2)^2} \qquad \begin{array}{l} R_{\text{KERNEL}} = 5000 \text{ AU} \\ \rho_{\text{KERNEL}} = 3 \times 10^{-18} \text{ g cm}^{-3} \\ R_{\text{CORE}} = 50\,000 \text{ AU} \end{array}$$

SPH code SEREN by David Hubber et al. (2011)

Radiative transfer method of Stamatellos et al. (2007)



lag column density

### The importance of radiative feedback from protostars

#### **Simulation II:**

Radiative feedback from protostars is fully included Continuous radiative feedback [similar to Offner, K<u>rumholz, Klein 2009 simulations]</u>

$$L_{\star} = \left(\frac{M_{\star}}{M_{\odot}}\right)^{3} L_{\odot} + f_{\rm rad} \frac{GM_{\star}\dot{M}_{\star}}{R_{\star}} \left(1 - \frac{R_{\star}}{2R_{\rm sink}}\right)$$

 $f_{\rm rad} = 0.75$ 

The fraction of the accretion energy that is radiated away



#### Accretion is episodic: FU Ori's

[see talks by Dunham, Hartmann, Offner]

#### FU Ori-type stars

Hartmann & Kenyon 1996, ARAA

rise time: 1-10 yr duration: 10s to a few 100s yr Accretion rate: a few  $10^{-4} M_{\odot}/yr$ Mass: 0.01-0.1  $M_{\odot}/event$ 



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### Accretion is episodic: Herbig-Haro objects

Episodic accretion onto a protostar results in episodic ejection of material.



Visible • WFPC2

Infrared • NICMOS

HH111 Hubble Space Telescope WFPC2 • NICMOS

Reipurth Nature 340, 42–45(1989)

NASA and B. Reipurth (CASA, University of Colorado) • STScI-PRC00-05

#### Accretion is episodic: the luminosity problem

 The luminosities of protostars are not high enough (Kenyon et al. 1990; Evans et al. 2009; Dunham et al. 2010)

$$0.5 \mathrm{M}_{\odot}/10^5 \mathrm{yr} \rightarrow \dot{M} = 5 \times 10^{-6} M_{\odot} \mathrm{yr}^{-1} \rightarrow L = \frac{GMM}{R_{\star}} \approx 25 L_{\odot}$$

• FU Ori type outbursts may happen for all protostars providing a solution to the luminosity problem :the luminosity is very high only during short events

#### The case for episodic accretion

- Thermal instability (Bell & Lin 1994)
- Binary companion (Bonnell & Bastien)
- Gravitational instabilities (Vorobyov & Basu 2005, 2006)
- Planet "blocking" (Lodato & Clarke 2004)

• Zhu, Hartmann et al. 2008-2010: The combined effect of different angular momentum transfer efficiencies of the gravitational instability (GI) and magneto-rotational instability (MRI).

#### GI: works better >10 AU from the star MRI: works better at <1 AU

MRI is initiated when  $T_M > 1400$  K in the inner disc region and the outburst starts. Stops when temperarture drops again.



## Episodic accretion: GI & MRI (e.g. Zhu et al. 2010)



# A phenomenological model of episodic accretion in hydrodynamic simulations



# A phenomenological model of episodic accretion (based on Zhu et al. 2010)

Stamatellos, Whitworth, Hubber, 2011, ApJ

Episodic accretion is initiated when

Λ

$$M_{\rm MRI} \simeq 0.13 \,{\rm M}_{\odot} \,\left(\frac{M_{\star}}{0.2 \,{\rm M}_{\odot}}\right)^{2/3} \,\left(\frac{\dot{M}_{\rm IAD}}{10^{-5} \,{\rm M}_{\odot} \,{\rm yr}^{-1}}\right)^{1/9}$$

 $M_{\rm IAD} \ge M_{\rm MRI} \ (T_M = 1400 \text{ K})$ 

Duration of an episodic  $\Delta t_{\rm MRI} \simeq 0.25 \,\rm kyr \, \left(\frac{\alpha_{\rm MRI}}{0.1}\right)^{-1} \, \left(\frac{M_{\rm MRI}}{0.13 \,\rm M_{\odot}}\right)$ 

Zhu et al. 2010a

$$\dot{M}_{\star,\text{EA}} = \frac{M_{\text{MRI}}}{\Delta t_{\text{MRI}}} e^{-\frac{(t-t_0)}{\Delta t_{\text{MRI}}}}, \ t_0 < t < t_0 + \Delta t_{\text{MRI}}$$

### The importance of radiative feedback from protostars

#### **Simulation III:**

**R**adiative feedback from protostars is fully included; the radiative feedback is not continuous but episodic



## Comparison Continuous vs episodic accretion/feeback





#### The role of episodic accretion in low-mass star formation

- duration of outburst ( $\Delta t_{MRI}$ )
- how often an outburst happens  $(T_{EA})$



 $\Delta t_{MRI} << T_{EA}$ 

 $T_{EA} > t_{dyn} \sim 10^3 yr$ 

#### **Duration of episodic accretion event**

$$\Delta t_{\rm MRI} \simeq 0.25 \,\rm kyr \, \left(\frac{\alpha_{\rm MRI}}{0.1}\right)^{-1} \, \left(\frac{M_{\rm MRI}}{0.13 \,\rm M_{\odot}}\right)$$

**Time interval between successive episodic accretion events** 

$$T_{\rm EA} \simeq 13 \, {\rm kyr} \, \left( \frac{M_{\star}}{0.2 \, {\rm M}_{\odot}} \right)^{2/3} \, \left( \frac{\dot{M}_{_{\rm IAD}}}{10^{-5} \, {\rm M}_{\odot} \, {\rm yr}^{-1}} \right)^{-8/8}$$

## **Observing fragmenting discs**

Stamatellos, Maury et al. 2011, MNRAS



## **Concluding remarks**

Episodic radiative feedback (due to episodic accretion) limits considerably the effects disc heating by the protostar

 Disc fragmentation is still possible; discs fragment to form low-mass stars, brown dwarfs and planetary-mass objects

The frequency of episodic accretion events may regulate low-mass star formation in different environments; where accretion events are more frequent (e.g. in high-density regions) fewer low-mass stars are expected.

## Different mechanisms may dominate in different regions?



## Observing the early stage of fragmenting discs is possible but improbable due to duration of the process

- $\Delta t_{\text{FRAGM}} = 1.5 \times 10^4 \text{ yr}$
- $\Delta t_{\text{CLASS 0}} = 1.5 \times 10^5 \text{ yr}$
- 10% of young protostars have such discs (enough to give 50% of the brown dwarf population)

~1% of Class 0 objects should have large unstable discs