

Max-Planck-Institut für Radioastronomie

Interactions between the SMBH SgrA\* and its Environment Central Massive Objects: The Stellar Nuclei – Black Hole Connection 22–25 June, ESO, Garching, Germany Andreas Eckart

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#### VLBI at 230 GHz (1.3 mm wavelength)





Yuan et al. 2009, Balbus & Hawley 1998, Balbus 2003



Yuan et al. 2009

Adiabatic Expansion of Source Components in the Temporary Accretion Disk of SgrA\*



Eckart et al. 2008, ESO Messenger Eckart et al. 2009, A&A 500, 935

# Pattern of a NIR spot orbiting at the ISCO



#### Pattern recognition against polarized red noise



total intensity

5σ

3σ

1σ

mean

polarization angle

polarization degree

Zamaninasab et al. 2009

# Pattern of a spot orbiting at the ISCO



Zamaninasab et al. 2009

# Pattern recognition against polarized red noise



Polarized flares as the signature of strong gravity are significant against randomly polarized red noise

Polarization data are consistent with the orbiting spot hypothesis

# NIR Polarized Flux Density from SgrA\*





Meyer, Eckart, Schödel, Duschl, Muzic, Dovciak, Karas 2006a Meyer, Schödel, Eckart, Karas, Dovciak, Duschl 2006b Eckart, Schödel, Meyer, Ott, Trippe, Genzel 2006

~4min prograde ~30min static for 3.6x10\*\*6Msol



## Cometary Sources: Shaped by a Wind from SgrA\*?



X7 polarized with 30% at PA -34+-10 Mie → bow-shock symmetry along PA 56+-10 includes direction towards SgrA\*

Muzic, Eckart et al. 2007, 2009, 2010

#### **Cometary Sources: Proper Motions**



Muzic, Eckart et al. 2010

## Cometary Sources: Shaped by a Wind from SgrA\*?



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#### **Cometary Sources: Source Location**

y (arcsec)

At least X7 is located within +-3.2" of the plane of the sky with a 67% probability.

-6 a 67% probability. x (arcsec) However, X3 may be co-spatial with the locatic of the mini-cavity (see y (arcsec) X7 also Zhao et al. 2009). Х3 mini-cavity -4 -6  $P(V > V_{PM}, r) = 1 - \frac{1}{\sigma^2} \int_0^{V_{PM}} v \, exp\left(\frac{-v^2}{2\sigma^2}\right) dv$ csec) Muzic, Eckart et al. 2010





optical depth throughout the shock

$$\tau(\lambda) = \tau_{abs}(\lambda) + \tau_{sca}(\lambda) = L \int_{a_{-}}^{a_{+}} n_d(a) C_{ext}(a,\lambda) da$$

extinction coefficient

$$C_{ext} = \pi a^2 (Q_{abs} + Q_{sca})$$



### Examined stellar wind sources

- •Late B-type main sequence stars (B7-8V)
- •Herbig Ae/Be stars
- •Central stars of Planetary nebulae (CSPN)
- •Low-luminosity Wolf Rayet (WR) stars (WC-type stars)
- •Main sequence stars
- •Dust-blob (X3, not X7)

Muzic, Eckart et al. 2010

Luminosity from emission and scattering

 $L = L_{sca} + L_{th}$  $L_{th} \propto B(T_d)(1 - e^{-\tau_{abs}})\epsilon_{th}$  $L_{sca} \propto d^{-2}\epsilon_{sca}P(\theta_{sca})e^{-\tau_{sca}}$ 

dust temperature close to the central star (VanBuren & McCray 1988; Krügel 2003)

$$T_d = 27 \, a_{\mu m}^{-1/6} \, L_{*,38}^{1/6} \, d_{pc}^{-1/3} \, \mathrm{K}$$

normalized scattering function

$$P(\theta_{sca}) = \frac{1 - g^2}{1 + g^2 - 2gcos(\theta_{sca})}$$





X7





X7

X3

#### Examined stellar wind sources

Not a wind from a single star of the GC young (He-)stars but only a collective wind from heavily mass loosing stars can potentially explain the bow-shock structure of X3 and X7.

However, such a global wind only emerges on scales of ~10", where as the distance of X7 and X3 is only 0.8" to 3.4" and the fact that the bowshocks are elognated and point towards SgrA\* is not explained.

**Standard Accretion Theory** 

With an estimated total wind mass loss - that could be accreted - of

$$\dot{M} \approx 2 \times 10^{-4} M_{\odot} y r^{-1}$$

that corresponds to an accretion rate in Eddington units of

$$m \approx 1 \times 10^{-3}$$

for a black hole mass of

$$(3-4) \times 10^6 M_{\rm O}$$

#### **Standard Accretion Theory**

From the bolometric luminosity of  $L \approx 2.1 \times 10^{36} erg/s$ 

one obtains an efficiency of

$$\eta_{eff} = \frac{L_{bol}}{M_{Edd} c^2} \approx 5 \times 10^{-6}$$

For the estimated accretion rate this would imply a much larger bolometric luminosity than actually observed:

$$L \approx 0.1 \times M_{Edd} c^2 \approx 4 \times 10^{43} erg/s$$

This is larger by a factor of about  $10^7$ 

SgrA\* is either very ineffciently accreting matter of very little of the matter available for accretion actually reaches the MBH, while the rest is blown away.

#### Dynamical Model for Accretion Ineraction between the 'starburst' and the black hole

Massloss from stars  $10^{-4} M_o yr^{-1}$ radiation efficiency of SgrA\*  $\approx 10^{-7}$ Bower et al. 2003: RM of linear polarized flux rules out large accretion rates

Quataert 2003: hydrodynamic calculations show that almost the entire mass gets blown away in a central wind. Available for accretion:  $< 10^{-5} M_{o} y r^{-1}$ 



## Cometary Sources: Shaped by a Wind from SgrA\*?



X7 polarized with 30% at PA -34+-10 Mie → bow-shock symmetry along PA 56+-10 includes direction towards SgrA\* Besides the Mini-Cavity – the strongest indication for a fast wind from SgrA\*!

Muzic, Eckart, Schödel et al. 2007, 2009, 2010

#### Sketch of an Outflow Model: The Combined Wind from the Cluster of Hot Stars and SgrA\*



Muzic, Eckart, Schödel et al. 2007 A&A 469, 993 Sabha, Eckart, Witzel, et al. 2009

#### GC SMBH interaction with the ISM

The following aspects need to be considered:

- accretion through a preferred plane?
- presence of a temporary relativistic disk
- inefficient accretion
- preferred direction of a wind from SgrA\* in addition to a wind from all mass loosing stars (min-cavity and bow-shock stars)

